

# MATH 201 LECTURE 28: STURM-LIOUVILLE THEORY

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- Many examples here are taken from the textbook. The first number in ( ) refers to the problem number in the UA Custom edition, the second number in ( ) refers to the problem number in the 8th edition.

## 0. REVIEW

- **Method of Separation of Variables.**

Given equation

$$\frac{\partial u}{\partial t} = \beta \frac{\partial^2 u}{\partial x^2} + P(x, t), \quad a < x < b; \quad u(x, 0) = f(x), \quad + \text{boundary conditions} \quad (1)$$

1. Require  $X(x)T(t)$  to solve the homogeneous equation

$$\frac{\partial u}{\partial t} = \beta \frac{\partial^2 u}{\partial x^2} \quad (2)$$

which leads to eigenvalue problem for  $X$ :

$$X'' - KX = 0 + \text{boundary conditions.} \quad (3)$$

Solve it to get  $X_n$  and  $K_n$ . Note that the natural range of  $n$  is **not** always  $1, 2, 3, \dots$

2. Expand

$$f(x) = \sum_n f_n X_n. \quad (4)$$

Expand

$$P(x, t) = \sum_n p_n(t) X_n. \quad (5)$$

3. Solve

$$T_n' - \beta K T_n = p_n(t), \quad T_n(0) = f_n \quad (6)$$

to obtain  $T_n$ .

4. Write down the solution

$$u(x, t) = \sum_n T_n(t) X_n(x). \quad (7)$$

- We have seen how this method works when  $f(x)$  and  $P(x, t)$  are already given in the form

$$f(x) = \sum_n f_n X_n; \quad P(x, t) = \sum_n p_n(t) X_n. \quad (8)$$

However in general this is not the case.

- Question: For arbitrary  $f(x)$ , is it possible to write it as  $f(x) = \sum_n f_n X_n$  with  $X_n$ 's the eigenfunctions obtained in Step 1? If so, how?
- Answer: Yes. See below.

## 1. STURM-LIOUVILLE THEORY

- **Sturm-Liouville theory**, developed almost 200 years ago by Jacques Charles François Sturm (1803 – 1855) and Joseph Liouville (1809 – 1882) studies the following problem: Given an general eigenvalue problem

$$-(p(x) X')' + q(x) X = \lambda w(x) X, \quad a < x < b \quad (9)$$

with boundary conditions

$$\alpha_1 X(a) + \beta_1 X'(a) = 0; \quad \alpha_2 X(b) + \beta_2 X'(b) = 0. \quad (10)$$

What can we say about the eigenvalues/eigenfunctions?

**Theorem 1. (Sturm-Liouville, Woolly version)** *The following hold true:*

1. *The eigenvalues are countable, and can be ordered by their sizes.*
2. *For each eigenvalue  $\lambda_n$ , the eigenfunction can be written as  $C X_n$ , where  $C$  is an arbitrary constant.*
3. *The  $X_n$ 's are "orthogonal" in the following sense:*

$$\int_a^b X_m(x) X_n(x) w(x) dx = 0 \text{ whenever } m \neq n. \quad (11)$$

4. *The  $X_n$ 's are "complete" in the following sense: Any reasonable  $f(x)$  (for example, bounded) has exactly one representation as linear combination of  $X_n$ 's:*

$$f(x) = \sum_n f_n X_n. \quad (12)$$

The "=" here means

$$\lim_{N \rightarrow \infty} \int \left| f(x) - \sum_{n < N} f_n X_n \right| dx = 0. \quad (13)$$

**Remark 2.** We intentionally choose not to present the precise version.

**Example 3.** Consider the eigenvalue problem

$$X'' - K X = 0; \quad X(0) = X(L) = 0; \quad (14)$$

We know that the eigenfunctions are

$$X_n = \sin\left(\frac{n\pi x}{L}\right), \quad n = 1, 2, 3, \dots \quad (15)$$

Then from the above theorem we know that any  $f(x)$  can be written as

$$f(x) = \sum_{n=1}^{\infty} f_n \sin\left(\frac{n\pi x}{L}\right). \quad (16)$$

We will see later that this expansion has a name: Fourier Sine Series.

**Example 4.** Consider the eigenvalue problem

$$X'' - K X = 0; \quad X'(0) = X'(L) = 0. \quad (17)$$

We know that the eigenfunctions are

$$X_n = \cos\left(\frac{n\pi x}{L}\right), \quad n = 0, 1, 2, 3, \dots \quad (18)$$

So the above theorem tells us any  $f(x)$  can be written as

$$f(x) = \sum_{n=0}^{\infty} f_n \cos\left(\frac{n\pi x}{L}\right). \quad (19)$$

Such expansion is called: Fourier Cosine Series.

**Remark 5.** It should be emphasized that, naturally, the  $f_n$ 's change when we pick a different set of  $X_n$ 's.

- How to compute  $f_n$ 's.

- Problem: Determine  $f_n$ 's in

$$f(x) = \sum_n f_n X_n. \quad (20)$$

- Idea: Use "orthogonality":

$$\int_a^b X_m(x) X_n(x) w(x) dx = 0 \text{ when } m \neq n. \quad (21)$$

- Let's set a particular  $n_0$  and try to find out  $f_{n_0}$ . As we try to use the above orthogonality, naturally we multiply both sides of

$$f(x) = \sum_n f_n X_n. \quad (22)$$

by  $X_{n_0}(x) w(x)$ , and then integrate from  $a$  to  $b$ . We have

$$\begin{aligned} \int_a^b f(x) X_{n_0}(x) w(x) dx &= \int_a^b \left[ \sum_n f_n X_n \right] X_{n_0}(x) w(x) dx \\ &= \sum_n f_n \int_a^b X_n(x) X_{n_0}(x) w(x) dx. \end{aligned} \quad (23)$$

As

$$\int_a^b X_n(x) X_{n_0}(x) w(x) dx = 0 \text{ for all } n \neq n_0 \quad (24)$$

we see that the right hand side in fact has exactly one nonzero term:

$$\int_a^b X_{n_0}(x)^2 w(x) dx. \quad (25)$$

Thus we reach

$$\int_a^b f(x) X_{n_0}(x) w(x) dx = f_{n_0} \int_a^b X_{n_0}(x)^2 w(x) dx. \quad (26)$$

and consequently

$$f_{n_0} = \frac{\int_a^b f(x) X_{n_0}(x) w(x) dx}{\int_a^b X_{n_0}(x)^2 w(x) dx}. \quad (27)$$

- Special cases most relevant to us:

- In all our problems the equation in the eigenvalue problem is

$$X'' - K X = 0 \implies w(x) = 1. \quad (28)$$

So the formula becomes

$$f_n = \frac{\int_0^L f(x) X_n(x) dx}{\int_0^L X_n(x)^2 dx}. \quad (29)$$

- Fourier Cosine and Fourier Sine Series.

Note that, as soon as we know  $X_n$ 's, the denominator  $\int_0^L X_n(x)^2 dx$  can be calculated beforehand, without knowledge of  $f(x)$ .

- Fourier Cosine Series.

In this case

$$X_n = \cos\left(\frac{n\pi x}{L}\right). \quad n = 0, 1, 2, 3, \dots \quad (30)$$

We have

$$\int_0^L \left[ \cos\left(\frac{n\pi x}{L}\right) \right]^2 dx = \int_0^L \frac{\cos\left(\frac{2n\pi x}{L}\right) + 1}{2} dx = \frac{L}{2}. \quad (31)$$

**Note that the above calculation is wrong when  $n = 0$ .** We have to calculate the  $n = 0$  case separately:

$$\int_0^L 1^2 dx = L. \quad (32)$$

So the  $f_n$ 's in the Fourier Cosine expansion

$$f(x) = \sum_{n=0}^{\infty} f_n \cos\left(\frac{n\pi x}{L}\right) \quad (33)$$

are given by

$$f_0 = \frac{1}{L} \int_0^L f(x) dx; \quad f_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx, \quad \text{for } n = 1, 2, 3, \dots \quad (34)$$

A more popular way of writing it is setting  $a_0 = 2 f_0$ , and  $a_n = f_n$  to get a universal formula

$$a_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx, \quad n = 0, 1, 2, 3, \dots \quad (35)$$

The Fourier cosine series then reads

$$\frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi x}{L}\right). \quad (36)$$

– Fourier Sine Series.

In this case

$$X_n = \sin\left(\frac{n\pi x}{L}\right), \quad n = 1, 2, 3, \dots \quad (37)$$

Similar calculation as in the previous case gives

$$f(x) = \sum_{n=1}^{\infty} f_n \sin\left(\frac{n\pi x}{L}\right) \implies f_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx, \quad n = 1, 2, 3, \dots \quad (38)$$

o Notation: Often, to emphasize the relation between Fourier Cosine/Sine series and Fourier series, the following notation is used:

– Fourier Cosine:

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi x}{L}\right), \quad a_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx, \quad n = 0, 1, 2, 3, \dots \quad (39)$$

– Fourier Sine:

$$f(x) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right), \quad b_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx, \quad n = 1, 2, 3, \dots \quad (40)$$

We will see in the next two lectures the reason of choosing the letters  $a, b$ .

## 2. EXAMPLES

• Fourier Cosine:

**Example 6. (10.4.13; 10.4 13)** Compute the Fourier cosine series for

$$f(x) = e^x, \quad 0 < x < 1. \quad (41)$$

**Solution.** We have  $T = 1$ . First

$$a_0 = \frac{2}{1} \int_0^1 e^x dx = 2(e - 1). \quad (42)$$

next

$$\begin{aligned}
 a_n &= 2 \int_0^1 e^x \cos(n\pi x) dx \\
 &= 2 \int_0^1 \cos(n\pi x) de^x \\
 &= 2 \left[ \cos(n\pi x) e^x \Big|_0^1 + n\pi \int_0^1 e^x \sin(n\pi x) dx \right] \\
 &= 2 [e(-1)^n - 1] + 2n\pi \int_0^1 \sin(n\pi x) de^x \\
 &= 2 [e(-1)^n - 1] + 2n\pi \left[ e^x \sin(n\pi x) \Big|_0^1 - n\pi \int_0^1 e^x \cos(n\pi x) dx \right] \\
 &= 2 [e(-1)^n - 1] - 2(n\pi)^2 \int_0^1 e^x \cos(n\pi x) dx \\
 &= 2 [e(-1)^n - 1] - (n\pi)^2 a_n.
 \end{aligned} \tag{43}$$

Therefore

$$a_n = \frac{2[e(-1)^n - 1]}{1 + (n\pi)^2}. \tag{44}$$

So the Fourier cosine series is given by

$$e^x = e - 1 + \sum_{n=1}^{\infty} \frac{2[e(-1)^n - 1]}{1 + (n\pi)^2} \cos(n\pi x). \tag{45}$$

- Fourier Sine:

**Example 7. (10.4.7; 10.4 7)** Compute the Fourier sine series for

$$f(x) = x^2, \quad 0 < x < \pi. \tag{46}$$

**Solution.** We have  $T = \pi$ . Compute

$$\begin{aligned}
 b_n &= \frac{2}{\pi} \int_0^{\pi} x^2 \sin(nx) dx \\
 &= -\frac{2}{n\pi} \int_0^{\pi} x^2 d\cos(nx) \\
 &= -\frac{2}{n\pi} \left[ x^2 \cos(nx) \Big|_0^{\pi} - 2 \int_0^{\pi} \cos(nx) x dx \right] \\
 &= -\frac{2}{n\pi} \left[ \pi^2 (-1)^n - \frac{2}{n} \int_0^{\pi} x d\sin(nx) \right] \\
 &= -\frac{2}{n\pi} \left[ \pi^2 (-1)^n - \frac{2}{n} \left( x \sin(nx) \Big|_0^{\pi} - \int_0^{\pi} \sin(nx) dx \right) \right] \\
 &= -\frac{2}{n\pi} \left[ \pi^2 (-1)^n - \frac{2}{n} \frac{1}{n} \cos(nx) \Big|_0^{\pi} \right] \\
 &= -\frac{2}{n\pi} \left[ \pi^2 (-1)^n - \frac{2}{n^2} [(-1)^n - 1] \right] \\
 &= \frac{2\pi}{n} (-1)^{n+1} + \frac{4}{n^3 \pi} [(-1)^n - 1].
 \end{aligned} \tag{47}$$

Therefore the Fourier sine series is

$$x^2 = \sum_{n=1}^{\infty} \left[ \frac{2\pi}{n} (-1)^{n+1} + \frac{4}{n^3 \pi} [(-1)^n - 1] \right] \sin(nx). \tag{48}$$

- More exotic examples.

**Example 8. (Mixed boundary conditions)** Expand  $f(x) = x$  into  $\sum f_n X_n$  with  $X_n$  eigenfunctions of

$$X'' - KX = 0, \quad 0 < x < \pi; \quad X(0) = X'(\pi) = 0. \tag{49}$$

**Solution.** We have already solved (in the previous lecture) the eigenfunctions:

$$X_n = \sin\left(\frac{2n+1}{2}x\right), \quad n = 0, 1, 2, 3, \dots \quad (50)$$

Now prepare:

$$\int_0^\pi X_n^2 = \int_0^\pi \frac{1 - \cos((2n+1)x)}{2} dx = \frac{\pi}{2} - \frac{1}{2(2n+1)} [\sin(2n+1)\pi - \cos 0] = \frac{\pi}{2} = \frac{\pi}{2}. \quad (51)$$

Thus

$$\begin{aligned} f_n &= \frac{2}{\pi} \int_0^\pi x \sin\left(\frac{2n+1}{2}x\right) dx \\ &= \frac{2}{\pi} \int_0^\pi x d\left[-\frac{2}{2n+1} \cos\left(\frac{2n+1}{2}x\right)\right] \\ &= -\frac{4}{\pi(2n+1)} \left[ \left(x \cos\left(\frac{2n+1}{2}x\right)\right) \Big|_0^\pi - \int_0^\pi \cos\left(\frac{2n+1}{2}x\right) dx \right] \\ &= -\frac{4}{\pi(2n+1)} \left[ \pi \cos\left(\frac{2n+1}{2}\pi\right) - \frac{2}{2n+1} \sin\left(\frac{2n+1}{2}x\right) \Big|_0^\pi \right] \\ &= -\frac{4}{\pi(2n+1)} \left[ \pi \cdot 0 - \frac{2}{2n+1} \left[ \sin\left(\frac{2n+1}{2}\pi\right) - 0 \right] \right] \\ &= \frac{8}{\pi(2n+1)^2} \sin\left(\frac{2n+1}{2}\pi\right) \\ &= \frac{8}{\pi(2n+1)^2} (-1)^n. \end{aligned} \quad (52)$$

So the expansion is

$$x = \sum_{n=0}^{\infty} \frac{8}{\pi(2n+1)^2} (-1)^n \sin\left(\frac{2n+1}{2}x\right). \quad (53)$$

### 3. POINTWISE CONVERGENCE?

- Recall that  $f(x) = \sum f_n X_n$  in the above means

$$\lim_{N \rightarrow \infty} \int \left| f(x) - \sum_{n < N} f_n X_n \right| dx = 0. \quad (54)$$

However in some applications we would like to know for a particular  $x_0$ , what is the value of

$$\left[ \sum_n f_n X_n \right](x_0) = \lim_{N \rightarrow \infty} \sum_{n < N} f_n X_n(x_0) \quad (55)$$

and

$$\lim_{N \rightarrow \infty} \int \left| f(x) - \sum_{n < N} f_n X_n \right| dx = 0 \text{ **does not imply** } \lim_{N \rightarrow \infty} \sum_{n < N} f_n X_n(x_0) = f(x_0). \quad (56)$$

Therefore we need to study “pointwise convergence” property of the expansions.

- For general  $X_n$  in the Sturm-Liouville Theorem, the situation seems quite complicated and I am yet to be sure of the existence of a complete theory.
- However, for the  $X_n$ 's obtained from the eigenvalue problem  $X'' - KX = 0$  + boundary conditions, we know exactly what  $\lim_{N \rightarrow \infty} \sum_{n < N} f_n X_n(x_0)$  is. In the next lecture we will reveal this through the study of Fourier series, which is the expansion of  $f(x)$  using  $\cos\left(\frac{2n\pi x}{L}\right)$  and  $\sin\left(\frac{2n\pi x}{L}\right)$ .
- Note that  $\cos\left(\frac{2n\pi x}{L}\right)$  and  $\sin\left(\frac{2n\pi x}{L}\right)$  are actually eigenfunctions to

$$X'' - KX = 0 \quad (57)$$

with periodic boundary condition (which is not included in the standard Sturm-Liouville theorem!),

## 4. ORTHONORMAL SET OF FUNCTIONS

- Any set of **nonzero** functions  $\{f_n\}_{n=1}^{\infty}$  satisfying, for some  $w \geq 0$

$$\int_a^b f_m f_n w \, dx = 0 \text{ whenever } n \neq m \quad (58)$$

is said to be an **orthogonal system** with respect to weight  $w$  on the interval  $[a, b]$ .

If furthermore we have

$$\int_a^b f_n^2 w \, dx = 1, \quad n = 1, 2, 3, \dots \quad (59)$$

then  $\{f_n\}$  is called an **orthonormal system**.

The main property of an orthogonal system is that if

$$f(x) \sim c_1 f_1 + c_2 f_2 + \dots \quad (60)$$

then the coefficients can be determined through

$$c_m = \frac{\int f(x) f_m(x) w(x) \, dx}{\int f_m^2(x) \, dx}. \quad (61)$$

If  $\{f_n\}$  is furthermore orthonormal, then

$$c_m = \int f(x) f_m(x) w(x) \, dx. \quad (62)$$

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**Example 9. (10.3.26, 10.3.27)** Show that the set of functions

$$\left\{ \cos \frac{\pi}{2} x, \sin \frac{\pi}{2} x, \dots, \cos \frac{(2n-1)\pi}{2} x, \sin \frac{(2n-1)\pi}{2} x, \dots \right\} \quad (63)$$

is an orthonormal system on  $[-1, 1]$  with respect to the weight function  $w(x) \equiv 1$ .

Then find the orthogonal expansion for

$$f(x) = \begin{cases} 0 & -1 < x < 0 \\ 1 & 0 < x < 1 \end{cases} \quad (64)$$

in terms of this orthonormal system.

**Solution.**

- Verify orthonormality.

1. Integrating product of different functions gives 0.

We compute for  $n \neq m$ ,  $n, m = 1, 2, \dots$  (note that  $w = 1$ )

$$\begin{aligned} \int_{-1}^1 \cos \frac{(2n-1)\pi x}{2} \cos \frac{(2m-1)\pi x}{2} \cdot 1 \, dx &= \frac{1}{2} \int_{-1}^1 \cos((n+m-1)\pi x) \, dx \\ &+ \frac{1}{2} \int_{-1}^1 \cos((n-m)\pi x) \, dx \end{aligned} \quad (65)$$

As neither  $n+m-1$  nor  $n-m$  is zero, we have

$$\int_{-1}^1 \cos((n+m-1)\pi x) \, dx = \frac{1}{(n+m-1)\pi} \sin((n+m-1)\pi x) \Big|_{-1}^1 = 0, \quad (66)$$

$$\int_{-1}^1 \cos((n-m)\pi x) \, dx = \frac{1}{(n-m)\pi} \sin((n-m)\pi x) \Big|_{-1}^1 = 0. \quad (67)$$

Similarly we compute

$$\int_{-1}^1 \sin \frac{(2n-1)\pi x}{2} \sin \frac{(2m-1)\pi x}{2} \cdot 1 \, dx = 0 \quad (68)$$

for  $n \neq m$ .

Finally we can compute

$$\int_{-1}^1 \cos \frac{(2n-1)\pi x}{2} \sin \frac{(2m-1)\pi x}{2} \cdot 1 \, dx = 0 \quad (69)$$

Note that this time  $n = m$  is OK.

2. Integrating the square of any function in the list gives 1.

We compute

$$\int_{-1}^1 \left( \cos \frac{(2n-1)\pi x}{2} \right)^2 \cdot 1 \, dx = \frac{1}{2} \int_{-1}^1 [1 + \cos((2n-1)\pi x)] \, dx = 1. \quad (70)$$

Similarly we have

$$\int_{-1}^1 \left( \sin \frac{(2n-1)\pi x}{2} \right)^2 \, dx = 1. \quad (71)$$

Thus the set of functions is an orthonormal system.

o Orthogonal expansion for

$$f(x) = \begin{cases} 0 & -1 < x < 0 \\ 1 & 0 < x < 1 \end{cases} \quad (72)$$

Recall that if

$$f(x) \sim c_1 f_1 + c_2 f_2 + \dots \quad (73)$$

then the coefficients can be determined through

$$c_m = \frac{\int f(x) f_m(x) w(x) \, dx}{\int f_m^2(x) \, dx}. \quad (74)$$

In case of our system we write

$$f(x) = \sum_{n=1}^{\infty} \left\{ a_n \cos \frac{(2n-1)\pi x}{2} + b_n \sin \frac{(2n-1)\pi x}{2} \right\} \quad (75)$$

and compute

$$\begin{aligned} a_n &= \int_{-1}^1 f(x) \cos \frac{(2n-1)\pi x}{2} \, dx \\ &= \int_0^1 \cos \frac{(2n-1)\pi x}{2} \, dx \\ &= \frac{2}{(2n-1)\pi} \sin \frac{(2n-1)\pi x}{2} \Big|_0^1 \\ &= \frac{2}{(2n-1)\pi} \sin(n\pi - \pi/2) \\ &= \frac{2(-1)^{n+1}}{(2n-1)\pi}. \end{aligned} \quad (76)$$

and

$$\begin{aligned} b_n &= \int_{-1}^1 f(x) \sin \frac{(2n-1)\pi x}{2} \, dx \\ &= \int_0^1 \sin \frac{(2n-1)\pi x}{2} \, dx \\ &= -\frac{2}{(2n-1)\pi} \cos \frac{(2n-1)\pi x}{2} \Big|_0^1 \\ &= \frac{2}{(2n-1)\pi}. \end{aligned} \quad (77)$$



Thus finally we have

$$f(x) \sim \sum_{n=1}^{\infty} \frac{2}{(2n-1)\pi} \left[ (-1)^{n+1} \cos \frac{(2n-1)\pi x}{2} + \sin \frac{(2n-1)\pi x}{2} \right]. \quad (78)$$

#### 5. NOTES AND COMMENTS

- Note that the boundary conditions

$$\alpha_1 y(0) + \beta_1 y'(0) = 0; \quad \alpha_2 y(L) + \beta_2 y'(L) = 0. \quad (79)$$

covers Dirichlet, Neumann, and Mixed boundary conditions, but not periodic boundary conditions.